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# AC Magnetic Field Tolerance for the HRS, and Shielding Possibilities

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**Abstract:**

# 1 Allowed AC field

The HRS will have maximum resolution with a slit size of 100 microns. Small magnetic fields orthogonal to the beam path will deflect the beam by distances of up to 1 millimetre, but these can be corrected easily if they are static. But if they are fluctuating for example because of a nearby wire carrying a 60 Hz current, however, they will compromise the resolution unless the net peak-to-peak deflection is below about 10 microns.

There are 3 vulnerable areas: (1) From object slit to the first dipole (0.8 m), (2) between the dipoles in the region of the multipole (1.6 m), (3) from second dipole to the mass slit. The most stringent requirement comes from region 2.

## 1.1 region 2

The optics from the multipole to the mass slit is a 90 degree phase advance: all accumulated angular deflections are converted to position deflection at the mass slit, and the conversion factor is 1.6 metre. Thus we find the allowable peak-to-peak deflection to be  $\Delta\theta = 10 \mu\text{m}/(1.6 \text{ m}) = 6.25 \mu\text{radians}$ .

The angle due to a field  $B$  is  $\Delta\theta = BL/(B\rho)$ , where  $L$  is the distance between the dipole. It happens that  $L = 1.6$  metres as well. Inverting, we get maximum allowable p-p field

$$B = B\rho \Delta\theta/L. \quad (1)$$

$B\rho$  is the magnetic rigidity of the beam  $= mv/q$ . The condition on  $B$  will be thus worst for the lightest beams, but on the other hand, the highest resolution is not needed for the very lightest beams. Referring to Marchetto's thesis, Fig. 1.13, we see that the lightest mass for which a resolution of 20,000 may be needed is about mass number 30. Assuming then this mass but with the highest energy of 60 keV, we find that  $B\rho = 2000$  Gauss-metres.

Plugging in the number, we find

$$B_{\text{p-p}} \leq 8 \text{ milliGauss} \quad (2)$$

## 1.2 regions 1&3

These regions are only half as long as region 2, and moreover the effect of an angular deflection declines linear from dipole to slit. This results in another factor of 2 for an overall factor of 4. Thus, for these regions, the maximum p-p tolerance is 30 milliGauss.

## 2 Allowed net current

From Ampere's law,  $\oint \vec{B} \cdot d\vec{l} = \mu_0 I$ , we find that the magnetic field is 2 milliGauss-metre per Ampere. Thus, an unbalanced net current of 1 Ampere at 1 metre causes 2 mG field. At 2 metres, this would be 1 mG per Amp. And so on. This is for currents in the plane of the separator. Current components orthogonal to this plane have much smaller effects.

In terms of peak-to-peak currents in the plane, the above results can be summarized as follows.

Region	Field	Current at 1 m	2 metres
2	8mG	4 Amps	8 Amps
1&3	32mG	16 Amps	32 Amps

## 3 Shielding

We can think of two cases: Placing AC field sources inside shielding, or placing shielding around the RIB ion beam. In addition, there are two effects to consider: eddy currents and ferromagnetic material. The first depends upon the conductivity of the shielding material, and the second depends upon the magnetic permeability.

### 3.1 Eddy Current Effect

Any metal can attenuate AC magnetic fields, but this effect is small at low frequency. The time scale  $\tau$  is given by [\[link\]](#)

$$\tau_e = \frac{\mu_0 \sigma t R}{1 + 1/\mu_r}, \quad (3)$$

where  $R$  is the cylinder radius,  $t$  is the shield thickness,  $\mu_r$  is the relative permeability,  $\sigma$  is the shield conductivity. As  $\mu_r \gg 1$ , we can drop the denominator. For steel, resistivity is  $1/\sigma \sim 16 \mu\Omega\text{cm}$ , so  $\mu_0 \sigma \sim 8/\text{m}^2$ . A shield like 2-inch EMT

conduit has  $t = 1.65$  mm,  $R = 25$  mm. Thus  $\tau_e \sim 0.3$  milliseconds. Converted to a frequency  $f_e = 1/(2\pi\tau)$ , this is  $\sim 500$  Hz. At 60 Hz, the effect will be minimal.

Notice that the effect depends on the cross-sectional area of the conduit, so how much material would be needed to create effective shielding? The 2-inch conduit has a mass of 2 kg per metre of length. To use the eddy current shielding effect, we would need  $f_e$  to be 100 times smaller, requiring conduit of 200 kg/m. This is unrealistic.

These considerations apply whether the shield is around the source or the beam.

## 3.2 Permeable material

### 3.2.1 shield around the current

This does not work. Ampere's law,  $\oint \vec{B} \cdot d\vec{l} = \mu_0 I$ , implies that a circular cylindrical shell of ferromagnetic material, carrying inside of it a current wire concentric with the shell, will have no effect on the magnetic field from that wire external to the shell. This is independent of the permeability of the shell. This can be seen as being due to the fact that the symmetry of the configuration forces the magnetic field to be a constant at a given radius and entirely in the azimuthal direction. If the symmetry is broken, for example if the wire is off-centre, the shield will re-distribute the fields somewhat, but the overall size will remain as it is for the symmetric case.

### 3.2.2 shield around the beam

An infinite cylindrical shield placed in and perpendicular to an ambient field will reduce this field by a factor

$$\eta = \frac{2R}{\mu_r t}, \quad (4)$$

The factor 2 would be 3/2 for a spherical geometry, and so is expected to vary for less regular geometries but remain of order 1.

Ordinary mild steel has permeability of only 1,000 or so and a remanent field of a few Gauss. Fields this large are undesirable even if DC. So mild steel is not considered in this application.

Mu-metal, or permalloy is better, having a permeability of up to  $10^5$  and remanent fields of a few mG. Commercially, this is available in many thicknesses. Let us assume result in fields of 80 mG and we wish a factor 10 reduction. Then for 1 metre size, we need no more than a foil as thin as  $50 \mu\text{m}$ . Since all the flux that would have been in the shielded region is now inside the foil, the field there is  $R/t$

times as strong;  $80 \text{ mG} \times 1 \text{ metre} / 0.05 \text{ mm} = 1600 \text{ Gauss}$ . This is still below saturation. The material is rather cheap though so we might use 10 times thicker sheet. It is still rather easy to bend and apply to cover the beam pipe, as we have done in the past for the ISIS optics box.

## 4 Conclusion

- Conductors carrying AC should be kept as far away from the particle beam as practical.
- Conductors carrying AC should be tightly paired to minimize unbalanced current.
- Wrapping conductors in ferromagnetic conduit is of little value.
- If resulting AC fields are still larger than  $8 \text{ mG}$ , mu-metal shielding should be applied to the HRS vacuum vessels (but not inside the dipoles). The thickness needed is less than  $1 \text{ mm}$ .