

# Shielding the Beam Pipe



Rick Baartman, TRIUMF

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# Basic E&M

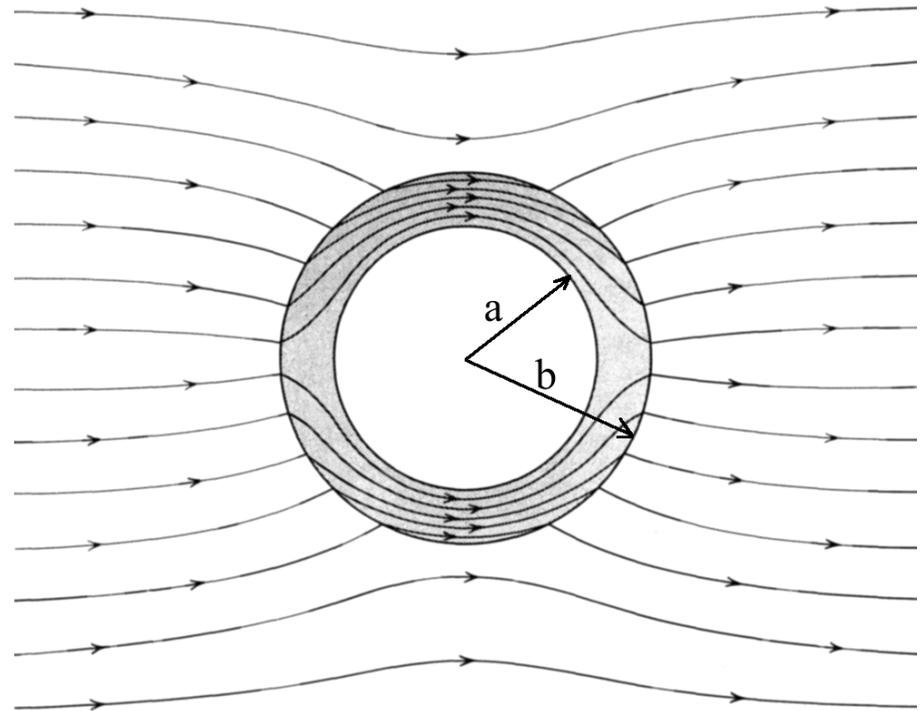
Familiar  
pic at right (borrowed from Jackson).

In polar  
coordinates  $(r, \theta)$ , with  $\theta = 0$  being  
the direction of the external field  
 $B_0$ , the field in the shielded region  
 $r < a$  is uniform and a constant:

$$B_i = \frac{4\mu}{(\mu + 1)^2 - (a/b)^2(\mu - 1)^2} B_0$$

(where  $\mu$  is the relative permeability) while the maximum field in the shield  
material is

$$B_\mu = \mu B_i$$



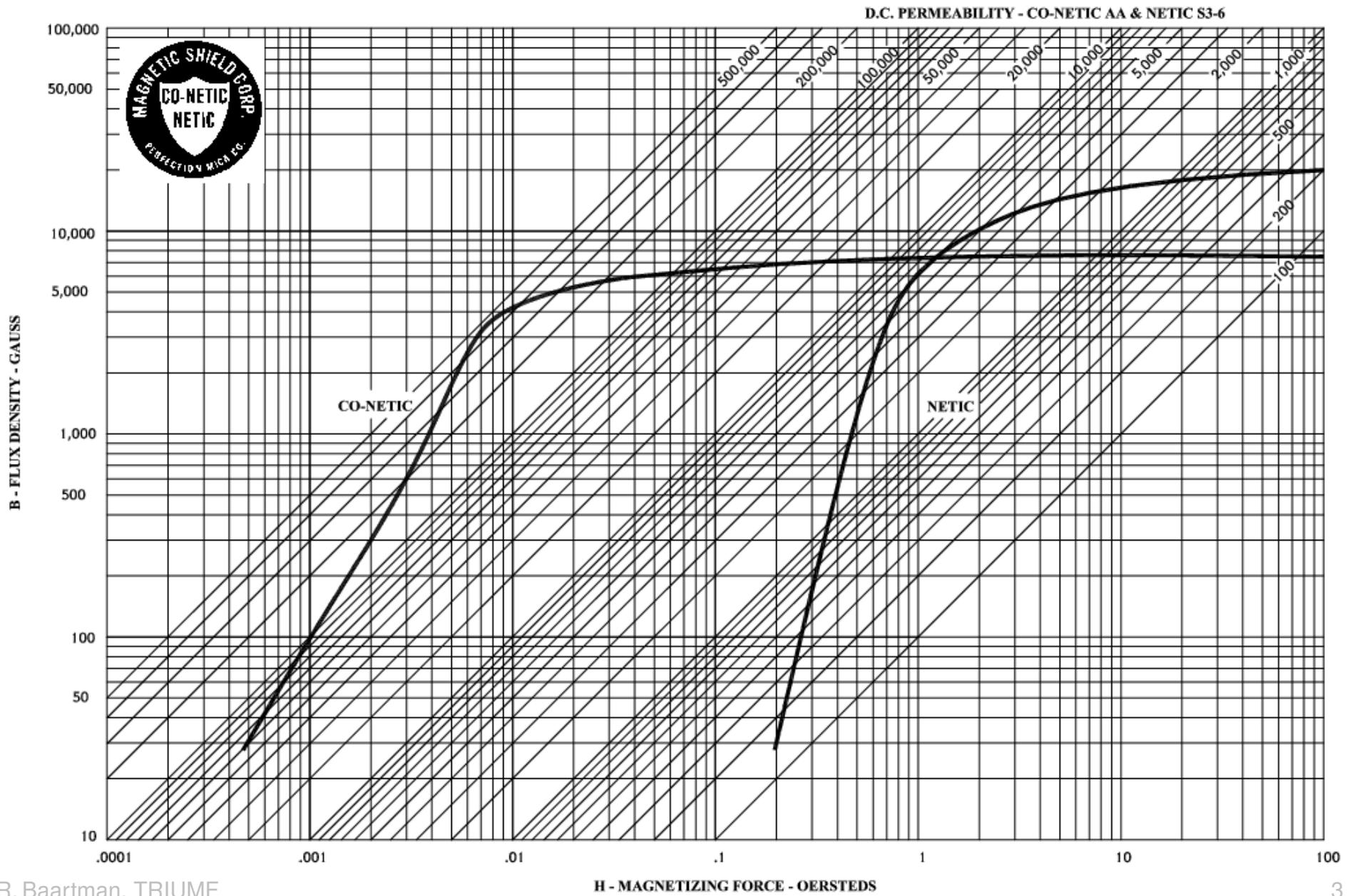
Let  $t$  be the thickness  $b - a$ . Then in the limits  $t \ll b$ ,  $\mu \gg 1$ , we have much simpler formulas:

$$B_i = \frac{2}{\mu} \frac{b}{t} B_0 \quad \text{and} \quad B_\mu = 2 \frac{b}{t} B_0$$

Roughly, the ambient field at worst is  $B_0 \sim 3$  G. As originally envisaged, we want  $B_i < 0.1$  G. We choose the standard material (commonly referred to as “Mu-metal”), which has  $\mu \sim 10^5$ , see “CO-NETIC” in figure below. (Co-netic AA is the trade name used by Perfection Mica Co. for Supermalloy.) The permeability is ill-defined as, owing to the irreversibility,  $B$  and  $H$  are not proportional.

This is available in foil with adhesive backing. ([Catalog here.](#)) Let us choose  $t = 0.10$  mm, since it is easy to handle and contributes little to the beam pipe size. We know  $b = 26$  mm, so

$$B_i = 16 \text{ mG}, \text{ and } B_\mu = 1.6 \text{ kG}.$$



Notice that the 1.6 kG field in the material is independent of  $\mu$ . Happily, this is well below saturation.

Looking at the permeability curves, we see that  $\mu$  is more nearly  $3 \times 10^5$ , and so  $B_i \sim 5$  mG. Further, the coercive force appears to imply fields less than a mG, so I conclude that coercivity will not contribute appreciably. (The manufacturer provides a [calculator](#), and this finds for 2-inch dia. shield in 3 G field:  $B_i = 4.3$  mG, in agreement.)

## Near the Quadrupoles

What will be the effect near quadrupoles?

It would be preferred to keep the shielding material well away from the quadrupole field, since as little as 10 G will saturate the shield.

Scaling from the 12Q12 field survey results, the fringe field would be 1/10 the peak field at a distance 47 mm from the quad centre, and dropping a factor of 2 every 10 mm. Most quads are at integrated strength of  $\sim 0.2$  T, with the peak field being  $\sim 0.1$  T. To drop a factor 100 to 10 G, requires a distance of 77 mm, meaning that the shield interruption is 154 mm in length. As the quad spacing is 2 m, the ambient field (3 G) is on average

$$\overline{B} = \frac{154 \text{ mm} \times 3 \text{ G} + (2000 - 154) \text{ mm} \times 0.0043 \text{ G}}{2 \text{ m}} = 0.23 \text{ G},$$

which is larger than the desired 0.1 G.

# Compensation

A possible solution is to compensate the field in the non-shielded regions. Simple Helmholtz coil loops can be used. As the quads are  $R_w = 224/\sqrt{2}$  mm axis to edge distance, the Helmholtz radius  $R$  that fits over the quad to give the necessary layout (coil loop separation =  $R$ ) is given by:

$$R = \frac{2}{3}R_w = 106 \text{ mm},$$

but we inflate this by 15 mm to allow for wire thickness, etc.

The well-known field formula

$$B = \frac{\mu_0 I}{R} \left( \frac{4}{5} \right)^{3/2}$$

gives

$$I = 40 \text{ A}.$$

# Compensation/Steering

This could be achieved with a pair of single loops, or divided in 10-100 turns to optimize power supply cost.

We could run many in series in regions where the ambient field is constant, or set individual currents, and this may obviate the need for some of the Radiabeam correctors.

They will certainly be cheaper and more convenient than those correctors, plus this would free up some space along the beam.

